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5	A Modified Method for Inferring Upper Cloud Top Height Using the GOES-12 Imager 10.7- and 13.3-µm Data
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25 Abstract

26	Accurate inference of upper transmissive cloud top height (CTH) is often difficult
27	using passive meteorological satellite instruments. Satellite retrievals using conventional CO ₂ -
28	slicing or CO ₂ -absorption techniques tend to systematically underestimate the upper CTHs.
29	These techniques are based on single-layer assumptions that the upper cloud occupies a
30	geometrically thin layer above a cloud-free surface. This study presents a new modified CO ₂ -
31	absorption technique (MCO2AT) to improve the inference of upper CTH. The MCO2AT
32	consists of three steps starting with a single-layer CO ₂ -absorption technique (SCO2AT)
33	followed by iterative retrieval procedures to obtain an enhanced upper CTH based on inferred
34	effective background radiances. The MCO2AT is applied to one-month of half-hourly
35	Geostationary Operational Environmental Satellite (GOES) imager 10.7- and 13.3-µm data.
36	Both the MCO2AT- and SCO2AT-inferred upper CTHs are compared with those obtained
37	from the ground-based Active Remotely Sensed Cloud Location (ARSCL) dataset taken at the
38	Atmospheric Radiation Measurement (ARM) Southern Great Plains Central Facility. For
39	overcast upper-cloud scenes, the mean upper CTH differences are $-0.33~(\pm 1.54)~\text{km}$ for
40	MCO2AT minus ARSCL in contrast to −1.58 (±2.08) km for SCO2AT minus ARSCL. For
41	broken upper-cloud scenes, the respective corresponding differences are $-2.66~(\pm 1.80)~\mathrm{km}$
42	and $-4.73~(\pm 2.28)$ km. Overall, the MCO2AT worked more effectively on multilayer clouds
43	and optically thin, geometrically thick clouds. The comparisons also revealed that more than
44	10% of the upper thin cirrus cases detected by the ARSCL were not retrieved by the
45	MCO2AT and SCO2AT, suggesting that a more sensitive cirrus cloud retrieval algorithm is
46	needed.

1. Introduction

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To infer upper cloud top height (CTH), current meteorological satellites usually employ a window technique that is based on analyses of the infrared (IR) and visible radiances obtained at atmospheric window channels [e.g., Rossow and Schiffer, 1991, 1999; Minnis et al., 1993, 2009] or a CO₂-slicing technique that is based on analyses of a suit of the multispectral IR sounding radiances obtained at various carbon dioxide (CO₂) absorption channels having nominal wavelengths of 13.3, 13.6, 13.9 and 14.3 µm [e.g., Chahine, 1974; McCleese and Wilson, 1976; Smith and Platt, 1978; Menzel et al., 1983, 1992; Smith and Frey, 1990]. Standard window techniques tend to underestimate the occurrence of highly transmissive cirrus clouds [Rossow, 1989; Wylie and Menzel, 1989, 1999; Jin et al., 1996]. Many of those missed clouds can be detected using the CO₂-slicing technique, which exploits the multiple IR sounding channels, with their varying spectral absorptivities and sensitivities, to infer the cloud-top pressure (CTP) and, hence, the CTH. Window techniques have typically been used to analyze 4-km data from the imagers on Geostationary Operational Environmental Satellite (GOES) I series, GOES-8 through GOES-11 [e.g., Smith et al., 2008], which measure radiances at five nominal wavelengths, 0.65, 3.9, 6.7, 10.7 and 12 µm [Menzel and Purdom, 1994]. The CO₂-slicing technique, on the other hand, has generally been applied to the relatively low spatial resolution (~12 km) GOES sounder data, which are taken at one visible and several IR channels including various CO₂-absorption bands [Menzel and Purdom, 1994; Schreiner et al., 2001]. However, beginning with GOES-12, the GOES-I imagers have a new 13.3-µm channel in place of the 12-µm channel. This channel modification was made to improve the cloud products derived from the GOES-I imagers [Schmit et al., 2001].

To that end, this study explores a new modified CO₂-absorption technique that employs the limited 10.7- and 13.3-µm channel pair for improving the inference of upper CTH. It has been shown that the performance of a CO₂-slicing technique depends on the channels used; different channel selections can result in different CTH solutions [Wielicki and Coakley, 1981; Eyre and Menzel, 1989]. In their theoretical study, Wielicki and Coakley [1981] noted that using the lower-wavelength sounding channels, which have the least atmospheric absorption, would retrieve the majority clouds throughout the troposphere. The higher sounding channels, which are more absorbing, would be limited mostly to detecting high clouds due to increased atmospheric opacity and reduced signal-to-noise ratios. However, it was also noted that the retrieval skill using the lower sounding channels would be degraded when the clear-sky radiances have significant uncertainties. While the uncertainties in clear-sky radiances generally imply the assumption of a single cloud layer, in more complex multilayer clouds the uncertainties also imply the impact by the presence of lower clouds underneath an upper cloud. In multilayer cases, the upwelling background radiances are no longer represented by the clear-sky radiances and the assumption of single-layer cloud can cause underestimations in the upper cloud CTHs [e.g., Chang and Li, 2005]. Among past validation studies, Wylie and Menzel [1989] compared the CO₂-slicing CTPs derived using the GOES VISSR Atmospheric Sounder (VAS) data to those derived from the ground-based lidar data, radiosonde moisture profiles, and GOES parallax measurements. Their results, which excluded multilayer and broken clouds, optically thin cirrus, and low-level stratocumulus, show that the averaged VAS-derived CTPs were within approximately ±50 hPa of the other CTP measurements. Frey et al. [1999] compared the CO₂-slicing CTHs derived using the Moderate-resolution Imaging Spectroradiometer

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(MODIS) Airborne Simulator (MAS) data to collocated airborne lidar data and found that the two CTHs agreed to within ±0.5 km for 32% and within ±1.5 km for 64% of the cases studied. Schreiner et al. [2001] compared CO₂-slicing CTHs derived from GOES-8 sounder data to the altitudes of apparent single-layer clouds estimated by Pilot Reports (PIREPs). They found that for high clouds the GOES-8 CTHs were, on average, ~910 m lower than the PIREPs data. Hawkinson et al. [2005] compared GOES-8 sounder CO₂-slicing retrievals of CTH to the ground-based Active Remotely Sensed Cloud Location (ARSCL) products [Clothiaux et al., 2000] of the Atmospheric Radiation Measurement (ARM) Program [Ackerman and Stokes, 2003], but they only used cases having relatively uniform GOES-8 CTHs. They found that the GOES-8 CTHs were lower than the ARSCL cloud tops by an average of 1.8 km with a standard deviation of 1.7 km.

More recently, Holz et al. [2006] evaluated the optically thin cirrus cases having a total optical depth less than 1 by comparing CTHs inferred from MAS data to collocated airborne Cloud Physics Lidar (CPL) data. They reported that for geometrically thick but optically thin clouds, the CO_2 -slicing CTHs were underestimated by more than 3 km. They suggest that the bias in CTH could be reduced using the hyper-spectral Scanning High-resolution Interferometer Sounder (S-HIS) measurements. Bedka et al. [2007] compared the GOES-12 sounder and imager data but studied only two cloud cases, one consisting of clouds at different levels and the other consisting entirely of single mid-level clouds obtained during the Atlantic-THORPEX Regional Campaign (ATReC). Comparisons with averaged airborne-CPL data show that for high clouds above 10 km, the CPL minus the GOES-12 sounder and imager CTHs yield mean biases of $\sim 3.8 \text{ km} \pm 2.9 \text{ km}$) and $\sim 3.1 \text{ km} \pm 2.7 \text{ km}$, respectively. For the other mid-level clouds, the mean biases were slightly less than 1 km (rms $\sim 1 \text{ km}$). In

comparisons with ARSCL data, Smith et al. [2008] found results similar to Hawkinson et al. [2005] using GOES-8 sounder retrievals of CTH for all single-layer clouds. However, for cirrus clouds having optical depths less than 3, the CO_2 -slicing results produced an average underestimate of 2.9 km \pm 2.0 km.

This study presents a new modified CO₂-absorbing technique aimed at improving on some of the common assumptions used in conventional single-layer CO₂-slicing techniques. The targeted assumptions are 1) the cloud occupies a single layer in the field of view (FOV) of the satellite instrument, 2) the single-layer cloud occupies infinitesimal or zero thickness by neglecting the cloud geometric thickness effect, 3) the emissivity of the cloud is the same in all sounding channels, and 4) the emissivity of the clear-sky surface is unity. Since conventional single-layer CO₂-slicing techniques systematically underestimate upper CTHs, the primary objectives of the modified CO₂-absorbing technique are: 1) to remedy the underestimation of upper cirrus CTH caused by the potential impact of underlying lower cloud, 2) to remedy the underestimation of single-layer cirrus CTH caused by the assumption of an infinitesimal thin cloud thickness, 3) to enhance the retrieval skill by reducing the impact of the uncertainties in surface emissivities and clear-sky radiances, and 4) to enhance the retrieval skill by relating the spectral cloud emissivities at the 10.7- and 13.3-μm channels.

In Section 2, our new modified CO₂-absorbing technique is delineated following a brief review of the conventional single-layer CO₂-slicing technique. Section 3 describes the data sets used herein. Section 4 presents comparisons of the results obtained from the conventional single-layer and our modified CO₂-absorbing techniques and examines those comparisons relative to the ground-based ARM ARSCL data. Section 5 gives the concluding remarks.

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2. Techniques

2.1. Single-layer CO₂-absorption Technique

- There are a handful of references on the CO₂-slicing or CO₂-absorption retrieval technique (see previous section). Since all operational CO₂-absorption techniques are based on a single-layer cloud assumption, hereafter, this method is referred to as the single-layer CO₂-absorption technique (SCO2AT). We briefly describe the SCO2AT that utilizes the radiance pair obtained by the GOES-12 imager: 10.7-μm (hereafter channel 4) and 13.3-μm channel (hereafter channel 6). [Note that the 12.0-μm channel is nominally channel 5, c.f. Schmit et al., 2001].
- Let us begin by considering a cloud-free pixel. The satellite-observed clear-sky radiance R_{clr}^{λ} at a spectral channel λ is given by

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$$R_{clr}^{\lambda} = B^{\lambda}(T_g)\xi^{\lambda}(P_g) + \int_{P_g}^{0} B^{\lambda}(T(P)) \frac{d\xi^{\lambda}(P)}{d\ln P} d\ln P, \qquad (1)$$

where $B^{\lambda}(T)$ denotes the Planck radiance at channel λ and temperature T, $\xi^{\lambda}(P)$ denotes the transmittance between atmospheric pressure level P and P = 0 (at the imager level), and the subscript g denotes the ground surface level. In (1), it is assumed that the surface emissivity = 1. Similarly, the opaque-cloud radiance R^{λ}_{ovc} observed for a completely overcast pixel at channel λ is given by

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$$R_{ovc}^{\lambda} = B^{\lambda}(T_c)\xi^{\lambda}(P_c) + \int_{P_c}^{0} B^{\lambda}(T(P)) \frac{d\xi^{\lambda}(P)}{d\ln P} d\ln P, \qquad (2)$$

where the subscript c denotes the cloud top level.

In the CO₂-IR channels, cloud reflectivity is assumed to be negligible and, thus, cloud transmissivity t_c^{λ} can be related to cloud emissivity e_c^{λ} by $t_c^{\lambda} = 1 - e_c^{\lambda}$. As such, the spectral radiance R_{pix}^{λ} for a cloudy pixel is given by

$$R_{pix}^{\lambda} = \varepsilon_c^{\lambda} R_{ovc}^{\lambda} + (1 - \varepsilon_c^{\lambda}) R_{clr}^{\lambda}, \tag{3}$$

- where $\varepsilon_c^{\lambda} = e_c^{\lambda} A_c$ is an effective cloud emissivity with A_c being the cloud cover fraction of the pixel. As a result, for opaque clouds (i.e., $e_c^{\lambda} = 1$), ε_c^{λ} is equivalent to the fractional coverage of the opaque cloud (i.e., A_c) and for completely overcast pixels (i.e., $A_c = 1$), ε_c^{λ} is equivalent to the emissivity of the overcast cloud (i.e., e_c^{λ}).
- Let us use the superscripts Ch4 and Ch6 for the 10.7-μm and 13.3-μm channels,
 respectively, and write (3) separately for the two channels as

$$R_{pix}^{\text{Ch4}} = \varepsilon_c^{\text{Ch4}} R_{ovc}^{\text{Ch4}} + (1 - \varepsilon_c^{\text{Ch4}}) R_{clr}^{\text{Ch4}}$$

$$\tag{4}$$

$$R_{pix}^{\text{Ch6}} = \varepsilon_c^{\text{Ch6}} R_{ovc}^{\text{Ch6}} + (1 - \varepsilon_c^{\text{Ch6}}) R_{clr}^{\text{Ch6}}. \tag{5}$$

We can manipulate (4) and (5) to derive the ratios of the two channels for the cloud radiative effects to yield

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$$\frac{R_{pix}^{\text{Ch6}} - R_{clr}^{\text{Ch6}}}{R_{pix}^{\text{Ch4}} - R_{clr}^{\text{Ch4}}} = \frac{\varepsilon_c^{\text{Ch6}} (R_{ovc}^{\text{Ch6}} - R_{clr}^{\text{Ch6}})}{\varepsilon_c^{\text{Ch4}} (R_{ovc}^{\text{Ch4}} - R_{clr}^{\text{Ch4}})}.$$
 (6)

- This two-channel approach obtained in (6) has been referred to as the radiance ratioing method in previous studies [Wielicki and Coakley, 1981; Eyre and Menzel, 1989].
- Previous applications of the radiance ratioing method in (6) assume that the spectral cloud emissivities at two different channels are the same (i.e., $\varepsilon_c^{\text{Ch4}} = \varepsilon_c^{\text{Ch6}}$). This assumption requires the two channels to be spectrally close. To infer the CTH, the left-hand side of (6)

represents the satellite-observed cloud radiative difference ratio (i.e., R_{pix}^{Ch4} and R_{pix}^{Ch6}) and the right-hand side of (6) represents the same ratio estimated using radiative transfer calculations (i.e., R_{ovc}^{Ch4} and R_{ovc}^{Ch6}). In practice, solution of the inferred CTH for a cloudy pixel is obtained by finding the optimal values of P_c and T_c that best satisfy (6) with a minimum difference between the two sides of the equation.

As seen in (6), however, the SCO2AT invokes the assumption of clear-sky background radiances (i.e., R_{clr}^{Ch4} and R_{clr}^{Ch6}). Since clear-sky radiances can not be directly observed in a cloudy pixel, they have to be obtained either from simulations based on radiative transfer calculations or from interpolations based on other cloud-free candidates found elsewhere. To this end, simulations based on radiative transfer calculations often serve as a better means because synoptic clouds in nature often extend over a large geographic area and for a long period of time, which hampers the measurement of nearby clear-sky radiances. The input data required for the radiative transfer calculations, atmospheric temperature and moisture profiles and surface temperatures and emissivities, are usually based on numerical weather analyses or forecasts. Nevertheless, Eq. (6) is not appropriate for a multilayer cloud situation.

2.2. Modified CO₂-absorption Technique

The modified CO₂-absorption technique (MCO2AT) is proposed to use inferred effective background radiances to replace the clear-sky radiances. In an idealized two-layer cloud situation, the effective background radiances represent the blackbody lower-cloud radiances. In reality, the effective background radiances are often a mixture of clear-sky radiances and/or lower-cloud radiances in single-layer and multilayer cloud situations.

Here let us consider an imager pixel containing two idealized cloud layers. We assume that the lower cloud layer can be represented by a slab of opaque cloud at an effective background level defined by the pressure P_{ebg} and temperature T_{ebg} , where the subscript ebg denotes the effective background. Following Eqs. (1) and (2) in the previous section, the effective background radiances R_{ebg}^{Ch4} and R_{ebg}^{Ch6} for channels 4 and 6, respectively, are given by

$$R_{ebg}^{Ch4} = B^{Ch4}(T_{ebg})\xi^{Ch4}(P_{ebg}) + \int_{P_{ebg}}^{0} B^{Ch4}(T(P)) \frac{d\xi^{Ch4}(P)}{d\ln P} d\ln P$$
 (7)

$$R_{ebg}^{\text{Ch6}} = B^{\text{Ch6}}(T_{ebg})\xi^{\text{Ch6}}(P_{ebg}) + \int_{P_{ebg}}^{0} B^{\text{Ch6}}(T(P)) \frac{d\xi^{\text{Ch6}}(P)}{d\ln P} d\ln P$$
 (8)

Then following equations (4)-(6), we obtain a new set of equations (9)-(11):

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$$R_{pix}^{\text{Ch4}} = \varepsilon_c^{\text{Ch4}} R_{ovc}^{\text{Ch4}} + (1 - \varepsilon_c^{\text{Ch4}}) R_{ebg}^{\text{Ch4}},$$
 (9)

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$$R_{pix}^{\text{Ch6}} = \varepsilon_c^{\text{Ch6}} R_{ovc}^{\text{Ch6}} + (1 - \varepsilon_c^{\text{Ch6}}) R_{ebg}^{\text{Ch6}}, \tag{10}$$

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$$\frac{R_{pix}^{\text{Ch6}} - R_{ebg}^{\text{Ch6}}}{R_{pix}^{\text{Ch4}} - R_{ebg}^{\text{Ch4}}} = \frac{\varepsilon_c^{\text{Ch6}} (R_{ovc}^{\text{Ch6}} - R_{ebg}^{\text{Ch6}})}{\varepsilon_c^{\text{Ch4}} (R_{ovc}^{\text{Ch4}} - R_{ebg}^{\text{Ch4}})}.$$
(11)

For different spectral cloud emissivities, we can relate $\varepsilon_c^{\text{Ch4}}$ and $\varepsilon_c^{\text{Ch6}}$ by

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$$\frac{\varepsilon_c^{\text{Ch6}}}{\varepsilon_c^{\text{Ch4}}} = \frac{1 - \exp(-\tau_c^{\text{Ch6}} / \mu)}{1 - \exp(-\tau_c^{\text{Ch4}} / \mu)},$$
 (12)

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$$\frac{\tau_c^{\text{Ch6}}}{\tau_c^{\text{Ch4}}} = \frac{\sigma_{ext}^{\text{Ch6}}}{\sigma_{ext}^{\text{Ch4}}},\tag{13}$$

- where τ_c^{Ch4} and τ_c^{Ch6} denote the corresponding spectral cloud optical depths, $\sigma_{ext}^{\text{Ch4}}$ and $\sigma_{ext}^{\text{Ch6}}$ denote the corresponding spectral extinction coefficients, and μ denotes the cosine of
- satellite viewing zenith angle.
- The modified Eq. (11) invokes the assumption of effective background radiances,
- R_{ebg}^{Ch4} and R_{ebg}^{Ch6} , which replace their counterparts in (6), the clear-sky radiances, R_{clr}^{Ch4} and
- R_{clr}^{Ch6} . The effective background radiances also can not be directly observed because the
- 225 upper cloud layer obscures the imager FOV. In practice, for the pixel-retrieval applications,
- we have developed an iterative retrieval procedure that applies the MCO2AT to infer
- 227 alternatively the effective background radiances and enhanced overcast radiances, R_{ovc}^{Ch4} and
- $R_{ovc}^{\rm Ch6}$, from which the effective background level $[P_{ebg}, T_{ebg}]$ and the improved upper CTH
- $[P_c, T_c]$ can be obtained. The iterative retrieval procedure consists of three main steps as
- described below. Step 1 basically applies the SCO2AT. If the SCO2AT retrieval in the first
- step is successful, the second step determines if the MCO2AT should be applied. The third
- step is then the essential component of the MCO2AT. The components of each step are listed
- 233 sequentially as follows.
- 1A. Initialize the effective background level at the ground level, i.e., $[P_{ebg}, T_{ebg}] = [P_g, T_g]$.
- 1B. Infer the clear-sky radiances R_{clr}^{Ch4} and R_{clr}^{Ch6} using Eq. (1).
- 1C. Retrieve an optimal solution set $[P_c, T_c]$ and R_{ovc}^{Ch4} and R_{ovc}^{Ch6} , which best satisfies Eq.
- 237 (6) using the single-layer assumption.
- 238 1D. If $P_c > 600 \text{ hPa or } P_c > (P_g 300 \text{hPa})$, go to Step 4C;
- Otherwise, proceed to Step 2A.

- 240 2A. Initialize $R_{ebg}^{Ch4} = R_{pix}^{Ch4}$ and use Eq. (7) to retrieve a rough solution set $[P_{ebg}, T_{ebg}]$.
- 241 2B. Use the rough solution set $[P_{ebg}, T_{ebg}]$ and Eq. (8) to infer a rough R_{ebg}^{Ch6}
- 242 2C. Compare the observed R_{pix}^{Ch6} with the rough R_{ebg}^{Ch6} .
- 243 2D. If $R_{pix}^{\text{Ch6}} R_{ebg}^{\text{Ch6}} > -\delta_R^{\text{Ch6}}$, go to Step 4B (here, δ_R^{Ch6} is the 13.3- μ m noise level);
- Otherwise, proceed to Step 3A.
- 245 3A. Take $[P_c, T_c]$ from Step 1C and take R_{ebg}^{Ch6} from Step 2B.
- 246 3B. Use $[P_c, T_c]$ and Eq. (2) for Channel 6 to infer R_{ovc}^{Ch6} , and then obtain $\varepsilon_c^{Ch6} =$
- 247 $(R_{pix}^{Ch6} R_{ebg}^{Ch6})/(R_{ovc}^{Ch6} R_{ebg}^{Ch6})$ from Eq. (10).
- 248 3C. Obtain $\varepsilon_c^{\text{Ch4}}$ through $\varepsilon_c^{\text{Ch6}}$ using Eqs. (11) and (12).
- 3D. Use $[P_c, T_c]$ and Eq. (2) for Channel 4 to infer R_{ovc}^{Ch4} , and then obtain R_{ebg}^{Ch4} =
- $(R_{pix}^{\text{Ch4}} \varepsilon_c^{\text{Ch4}} R_{ovc}^{\text{Ch4}}) / (1 \varepsilon_c^{\text{Ch4}}) \text{ from Eq. (9)}. \text{ Here, } R_{ebg}^{\text{Ch4}} \text{ is bounded by a maximum at}$
- 251 $R_{ebg}^{Ch4} = R_{clr}^{Ch4}$ and by a minimum at $R_{ebg}^{Ch4} = (R_{clr}^{Ch4} + R_{ovc}^{Ch4})/2$.
- 3E. Use the new R_{ebg}^{Ch4} and Eq. (7) to retrieve a new solution set $[P_{ebg}, T_{ebg}]$.
- 253 3F. Use the new solution set $[P_{ebg}, T_{ebg}]$ and Eq. (8) to infer new R_{ebg}^{Ch6} .
- 3G. Retrieve the optimal solution set $[P_c, T_c]$ and R_{ovc}^{Ch4} and R_{ovc}^{Ch6} , which best satisfies Eq.
- 255 (11) in the modified method.
- 3H. If $\left| \text{new } R_{ebg}^{\text{Ch6}} \text{old } R_{ebg}^{\text{Ch6}} \right| > \delta_R^{\text{Ch6}}$, repeat Steps 3B-3H; Otherwise, go to Step 4A.
- 257 4A. The solution set $[P_c, T_c]$ is retrieved from the MCO2AT with $[P_{ebg}, T_{ebg}]$. Stop.
- 258 4B. The solution set $[P_c, T_c]$ is retrieved from the SCO2AT with $[P_{ebg}, T_{ebg}] = [P_g, T_g]$.
- 259 Stop.

4C. No upper cloud (i.e., no $P_c < 600 \text{ hPa}$) is retrieved. Stop.

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The above retrieval steps will give a cloud pixel a possible upper CTH solution set $[P_c, T_c]$ inferred by the MCO2AT (4A) for an enhanced CO₂-CTH or by the SCO2AT (4B) for a regular CO₂-CTH. The upper CTH is inferred by finding the altitude corresponding to P_c in an appropriate vertical profile of pressure levels. The threshold of 600 hPa (~4.3 km) is selected conservatively because the CO₂-absorption method becomes less certain towards lower troposphere due to its reduced signal-to-noise ratio. There may be a possible low-level CTH solution (4C) having $P_c \ge 600$ hPa, which would require the window technique and/or a more robust cloud retrieval algorithm, but that is beyond the scope of this study. Here in this paper, we only focus on the results derived from the SCO2AT and MCO2AT. Note that from Step 1D, a SCO2AT-retrieved P_c is considered valid only when it is smaller than a threshold at 600 hPa or P_g –300hPa (the smaller of the two). From Step 2D, the MCO2AT is invoked only when the observed 13.3- μm radiance ($R_{\it pix}^{\rm Ch6}$) is smaller than the modeled 13.3- μm radiance (R_{ebg}^{Ch6}) by more than a threshold of $\delta_R^{Ch6} = 0.1~{\rm Wm}^{-2}{\rm sr}^{-1}{\rm \mu m}^{-1}$. This also means that, if $R_{pix}^{\text{Ch6}} > R_{ebg}^{\text{Ch6}} - \delta_R^{\text{Ch6}}$, the SCO2AT-retrieved P_c is adopted and no MCO2AT is invoked. The threshold $\delta_{\it R}^{\rm Ch6} = 0.1~{\rm Wm}^{-2}{\rm sr}^{-1}{\mu}{\rm m}^{-1}$ used here is slightly larger than five times the instrument noise level of the GOES-12 imager 13.3-µm channel. It is worth noting that the aforementioned iterative retrieval procedure is applicable to all single-layer and multilayer cloud pixels, except that the MCO2AT works more effectively for transmissive cirrus clouds in multilayer situations or geometrically thick, tenuous clouds. Figure 1 schematically illustrates the main differences between SCO2AT (Figs. 1a-1c, left side) and MCO2AT (Figs. 1d-1f, right side) and the various impacts of the effective

background radiances on three representative cases, namely 1) cirrus-overlying-stratus cloud (Figs. 1a and 1d), 2) geometrically thick, optically thin cloud (Figs. 1b and 1e), and 3) convective cloud (Figs. 1c and 1f). In the figure, all upper clouds are treated as a geometrically thin layer with uniform $[P_c, T_c]$. The SCO2AT cases on the left side depict the single-layer assumption with a fixed background level $[P_g, T_g]$ and the MCO2AT cases on the right side show the modified assumption with an enhanced background level $\left[P_{ebg},T_{ebg}\right]$. For the cirrus-overlying-stratus cloud situation (Figs. 1a and 1d), the effective background level is higher than the ground level (i.e., $P_{ebg} < P_{g}$) and the MCO2AT-inferred CTH (P_{c} in Fig. 1d) is raised higher than the SCO2AT-inferred CTH (P_c in 1a). For the geometrically thick, but optically thin cloud situation (Figs.1b and 1e), the effective background level is also higher than the ground level and likewise the MCO2AT-inferred CTH (P_c in Fig. 1e) is also higher than the SCO2AT-inferred CTH (P_c in Fig. 1b). However, for the third situation of a convective cloud (Figs. 1c and 1f), despite the effective background level being much higher than the ground level, the MCO2AT-inferred CTH (P_c in Fig. 1f) is only slightly higher than the SCO2AT-inferred CTH (P_c in Fig. 1c). This is because the convective cloud is more opaque than the upper transmissive clouds in the previous two situations. Hence, the upper CTHs inferred by both MCO2AT and SCO2AT are less sensitive to the different effective background radiances. We need to note that for opaque upper clouds, the MCO2AT retrieval procedure as in Step 3D prevents the effective background level from going too high in altitude. As such, a midway threshold chosen to be between the clear-sky and overcast radiances, i.e., $R_{ebg}^{\rm Ch4}=(R_{clr}^{\rm Ch4}+R_{ovc}^{\rm Ch4})/2$, is used to constrain a minimum P_{ebg} for the upper bound and a threshold at the maximum $P_{ebg} = P_g$,

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where $R_{ebg}^{\text{Ch4}} = R_{clr}^{\text{Ch4}}$ and surface emissivity = 1 are used to constrain the lower bound at the surface level. As a result, for opaque upper clouds similar CTHs are inferred from both the MCO2AT and SCO2AT.

3. Data

3.1. Model Data

The input atmospheric profiles of temperature, pressure, height, and moisture for calculating the clear-sky, overcast, and effective background radiances (R_{clr}^{λ} , R_{ovc}^{λ} and R_{ebg}^{λ}) are taken from the NOAA National Centers for Environmental Prediction (NCEP) 1-hourly Rapid Update Cycle (RUC) model-analysis dataset with a 40-km spatial resolution [Benjamin et al., 2004a, b]. The hourly RUC data cover a large portion of North America including the Contiguous United States (CONUS). The widely-used MODTRAN4 radiative transfer code [Berk et al., 1999] is used for simulations of the spectral radiances. For the spectral cloud emissivities at 10.7- and 13.3-µm channels, they are related through the adoption of an average extinction coefficient ratio of $\sigma_{ext}^{Ch6}/\sigma_{ext}^{Ch4} = 1.12$ (c.f. Eq. (13)) for an approximate range of $\sigma_{ext}^{Ch6}/\sigma_{ext}^{Ch4} = 1.02-1.25$ for various non-spherical ice crystal sizes [Yang et al., 2001]. This average extinction coefficient ratio results in an increase of approximately 10% in the calculated cloud emissivity ratio $\varepsilon_{ext}^{Ch6}/\varepsilon_{ext}^{Ch4}$ (c.f. Eq. (12)).

3.2. GOES-12 and ARSCL Data

To evaluate the performance of the MCO2AT-retrieved CTHs via comparisons with the ARSCL-derived CTHs, this study uses the half-hourly GOES-12 imagery data (scanning

at about 15 and 45 minutes of the UTC hour) obtained during the month of May 2005. During this month, the ARM field research location at the Southern Great Plains (SGP) Cloud and Radiation Test Bed (CART) site in Lamont, Oklahoma was frequently cloudy, having various cloud types almost every day. The half-hourly GOES-12 imagery data for the SGP domain were taken from the NASA Langley ARM imagery and cloud product archives [Ayers et al., 2006; see http://www-angler.larc.nasa.gov/]. Those images have a convolved 4 km \times 3.2 km spatial resolution at nadir. The original scanning resolution is about 4 km \times 2.3 km (north-south direction × east-west direction) for the 10.7-µm channel and about 8 km × 2.3 km (north-south × east-west) for the 13.3-µm channel. At the ARM SGP CART site [36.6°N, 97.5°W], the GOES-12 imager has an oblique satellite viewing zenith angle of ~48° and this degrades its spatial resolution by approximately a factor of 1.5 compared to the nadir. The ARSCL data products provide a time series of ground-based observations of cloud-top and cloud-base heights for single-layer or multilayer clouds observed at the ARM SGP CART site [Clothiaux et al., 2000]. The ARSCL algorithm combines the ground-based active remote sensing measurements from both a millimeter-wavelength cloud radar (MMCR) and a micro-pulse lidar (MPL) to estimate the locations of cloud-top and cloud-base heights at an up-looking spatial resolution of ~45 m and a temporal sampling rate at 10-second intervals. The ARSCL produces a much finer-resolution time series data than the half-hourly GOES-12 imager data. But at each observation time, the GOES-12 imager views a much larger cloud area than that viewed by the vertically pointing narrow-beamed ARSCL radar and lidar. Meticulous comparisons between the satellite- and ground-based cloud observations are needed to minimize the uncertainties in matching the data obtained from two different platforms.

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4. Comparisons

4.1. Comparisons of MCO2AT and SCO2AT CTPs

352 Figures 2 and 3 are used to demonstrate the SCO2AT and MCO2AT inferred CTPs. 353 Figure 2 shows an example of the GOES-12 imagery data for the 13.3-µm (2a) and 10.7-µm 354 (2b) channels obtained at 10:45 UTC on 1 May 2005 for SGP spatial domain, which extends 355 from 32°N to 42°N and from 105°W to 91°W. The ARM SGP CART site is indicated by the 356 little square near the center of the images. In Fig. 2a, the 13.3-µm brightness temperatures are 357 obscured by the CO₂ absorption and thus appear colder (brighter) than the 10.7-µm brightness temperatures in Fig. 2b. Figure 3 shows the corresponding SCO2AT-inferred P_c (Fig. 3a) 358 359 and P_g (Fig. 3b) and MCO2AT-inferred P_c (Fig. 3c) and P_{ebg} (Fig. 3d) for the images shown 360 in Fig. 2. The MCO2AT-inferred P_c (Fig. 3c) are generally smaller (higher in altitude) than the SCO2AT-inferred P_c in Fig. 3a. Similarly, the MCO2AT-inferred P_{ebg} values in Fig. 3d 361 362 are generally less than P_g in Fig. 3b. 363 In this example (Figs. 2 and 3), the upper clouds do not appear to be opaque as 364 evidenced by the relatively low MMCR reflectivities seen between 5 and 11 km in Figure 4, 365 which shows the corresponding MMCR reflectivity profiles obtained at the SGP CART site 366 on this day. Note that the MMCR reflectivity profiles are plotted in reverse UTC time so that 367 the reflectivity profiles for the time prior to 10:45 UTC (GOES-12 image time) are plotted 368 eastwards to correspond with the GOES-12 image clouds on the east side of the SGP CART 369 site, and likewise the reflectivity profiles for the time after 10:45 UTC are plotted westwards 370 to correspond with the GOES-12 image clouds on the west side of the SGP CART site.

Figure 4 shows that the clouds extend vertically several kilometers and the upper clouds are sometimes contiguous with but sometimes decoupled from the lower clouds. Such vertical extensions and split cloud layers make the satellite cloud retrieval difficult when the singlelayer and/or the geometrically thin layer assumptions are used. Both assumptions result in underestimations of satellite-inferred CTH; for example, the larger SCO2AT-inferred P_c by assuming a clear background level at P_g than the relatively smaller MCO2AT-inferred P_c . Figure 5 compares the differences between the MCO2AT- and SCO2AT-inferred P_c (5a) and the differences between their corresponding P_{ebg} and P_{g} (5b) for the cloud pixels shown in Fig. 3. In the figure, when both MCO2AT and SCO2AT have the same P_c , the retrieval Steps 2A-2D were used and their associated values of P_{ebg} and P_{g} are also identical. For pixels having P_c from MCO2AT that are smaller than their SCO2AT counterparts, retrieval Steps 3A-3H were invoked and the associated mean (standard deviation) P_c is about 381 (79) hPa for the MCO2AT and 451 (91) hPa for the SCO2AT. Their corresponding $P_{\rm ebg}$ and P_g means (standard deviation) are about 703 (61) hPa and 944 (43) hPa, respectively. The MCO2AT significantly decreased the upper-cloud P_c and the mean difference (standard deviation) of MCO2AT minus SCO2AT is about -70 (37) hPa for P_c and -241 (65) hPa for P_{ebg} minus P_{g} .

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Figure 6 shows GOES-12 13.3- μ m (Fig. 6a) and 10.7- μ m (Fig. 6b) images obtained at 10:45 UTC, 24 May 2005, when a deep convective storm system occupied a sizeable portion of the scene. Again, the associated SCO2AT-inferred P_c (Fig. 7a) and P_g (Fig. 7b) values exceed their MCO2AT-inferred P_c (Fig. 7c) and P_{ebg} (Fig. 7d) counterparts. Figure 8 shows

the corresponding SGP CART MMCR reflectivity profiles. Although such convective systems are generally thought to have more opaque cloud tops, the upper clouds as revealed in Fig. 8 have generally low radar reflectivities (negatives in dBZ) that do not appear as opaque as expected. The low or missing radar reflectivities are likely caused by attenuation by the lower clouds and by the inability of cloud radars to detect the small ice crystals and low ice water contents in the top portions of those clouds [e.g., McGill et al., 2004].

For the cloudy pixels in Figure 7, the mean difference (standard deviation) between the values of P_c from MCO2AT and SCO2AT is about -60 (40) hPa and the corresponding mean difference (standard deviation) of P_{ebg} minus P_g is about -318 (113) hPa. The differences between the MCO2AT- and SCO2AT-inferred P_c are substantial even around the central areas of the deep convective clouds (c.f. Figs. 7a and 7c).

4.2. Comparisons with ARSCL CTHs

In order to compare with the ARSCL CTH data, the MCO2AT- and SCO2AT-retrieved pixel-scale P_c are first converted to CTH z_c using the RUC vertical profiles of atmospheric pressure. Secondly, from each half-hourly GOES image analysis, spatial averages of z_c from the MCO2AT- and SCO2AT were computed based on the imager pixels obtained within a 15 km ×15 km area centered at the SGP CART site. Figure 9 shows the half-hourly mean z_c values from the MCO2AT (solid points) and SCO2AT (cross points) obtained on 1 May (9a) and 24 May (9b) 2005, along with the ARSCL time-series cloud vertical profile mask plotted in grey areas. Both Figs. 9a and 9b are also plotted in reverse

UTC time, in the same manner used for the MMCR reflectivity profiles in Fig. 4 and Fig. 8, to reflect the east-west cloud orientation in satellite images.

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The differences between the MCO2AT- and SCO2AT-inferred mean z_c is approximately 2 km at 10:45 UTC on May 1 (Fig. 9a) and is slightly less than 1 km at 10:45 UTC on May 24 (Fig. 9b). In the figure, the solid circles indicate that 100% of the pixels within the 15-km area had valid MCO2AT and SCO2AT retrievals whereas the solid triangles indicate that only a portion of the pixels in the 15-km area had valid MCO2AT and SCO2AT retrievals. (More details on the overcast and broken cloud scenes are presented later in Fig. 10.) The MCO2AT-inferred mean z_c values show generally good agreement with the ARSCL uppermost CTHs, whereas the SCO2AT-inferred mean z_c values are generally lower than the MCO2AT mean z_c by about 1-2 km. Some of the differences between the MCO2AT and SCO2AT mean z_c and the ARSCL data in Fig. 9b are larger than 3 km near the cloud edges or when the upper clouds are thin. To evaluate the performance of the MCO2AT relative to the SCO2AT, we analyzed a total of 1488 half-hourly data from the GOES-12 imager obtained during 1-31 May 2005 and examined both the MCO2AT and SCO2AT retrievals for the 15-km area at the SGP site. To facilitate the comparisons, the 1488 half-hourly 15-km scenes were divided into different categories according to both the imager-retrieved pixel z_c and the ARSCL uppermost CTH data within ± 1.5 minutes of the GOES-12 imager scan time at the SGP CART site

upper-cloud scene (solid circles in Fig. 9), requires that 100% of the GOES-12 imager pixels

(approximately at 17 and 47 minutes after the UTC hour). The first category, the overcast

of the ARSCL profile data within the ± 1.5 -min period of the imager scan-passing time must have a valid uppermost CTH greater than 4.3 km (~ 600 hPa) at every shot. The second category, the broken upper-cloud scene (solid triangles in Fig. 9), requires that either the GOES-12 imager pixels only have a portion with valid z_c retrievals in the 15-km area or the ARSCL profiles only have a portion with valid uppermost CTH greater than 4.3 km within the ± 1.5 -min period. The third category is called the border upper-cloud scene (as marked by open triangles in Fig. 11). In this category, the GOES-12 imager pixels have some valid z_c and the ARSCL profiles have no CTH > 4.3 km within the ± 1.5 -min period, but have some valid CTH > 4.3 km in a bigger window of ± 3 -min period.

Among the 1488 cases, we found 377 cases that were overcast upper-cloud scenes, 248 cases had broken upper-cloud scenes, 26 cases had border upper-cloud scenes, and 16 cases had missing data. There were also 196 cases having ARSCL profiles with valid CTH > 4.3 km within the ± 1.5 -min period but with no valid MCO2AT and SCO2AT z_c in the 15-km area. These cases are generally optically thin clouds with MMCR reflectivities less than -20 dBZ and, in many of the cases, it was necessary to widen the area to find valid z_c retrievals in the neighborhood. Those cases are not included in this study.

Figure 10 shows the one-to-one comparisons between the ARSCL and the MCO2AT (left column) and SCO2AT (right column) obtained mean CTHs for 1) overcast upper-cloud scenes (a and b), 2) broken upper-cloud scenes (c and d), and 3) border upper-cloud scenes (e and f). Here the mean ARSCL upper CTHs are obtained by averaging the 25th–75th percentiles of all valid uppermost CTHs > 4.3 km in the ± 1.5 -min period. The vertical bars in the figure indicate the range of the 25th–75th percentiles. The mean MCO2AT and SCO2AT

CTHs are obtained by averaging all valid pixel-retrieved z_c within ± 3 standard deviations of the mean value. For reference, the solid circles in Figs. 10a and 10b denote the overcast upper-cloud scenes obtained on both 1 May (Fig. 9a) and 24 May (Fig. 9b) and the solid triangles in Figs. 10c and 10d denote the broken upper-cloud scenes also obtained on the two days.

The results obtained here show significant improvements in CTH from the MCO2AT compared to the SCO2AT for overcast and broken upper-cloud scenes. Mean biases of the MCO2AT minus ARSCL CTHs are $-0.33~\text{km}\pm1.54~\text{km}$ for the overcast (Fig. 10a) and $-2.66~\text{km}\pm1.80~\text{km}$ for the broken scenes (10c). For SCO2AT minus ARSCL, the mean biases in CTH are $-1.58~\text{km}\pm2.08~\text{km}$ for the overcast (10b) and $-4.73~\text{km}\pm2.28~\text{km}$ for the broken scenes (10d). From Figs. 10a and 10b, it is worth mentioning that many MCO2AT- and SCO2AT-inferred CTHs are significantly larger than the ARSCL uppermost CTHs. We examined the MMCR reflectivities for these cases and found that they are mostly from precipitating convective clouds as revealed by relatively large dBZ values. Many of these cases are found on May 13, 23, 26, and 31 and the ARSCL uppermost CTHs were considerably lower than the MCO2AT- and SCO2AT-inferred mean CTHs.

Figure 11 shows the results for May 26. Between about 10:00 and 12:00 UTC, the ARSCL uppermost CTHs are systematically several kilometers lower than the MCO2AT and SCO2AT overcast upper CTHs. Most likely, the differences are due to attenuation of the cloud radar and lidar beams and the limitations of the radar for detecting small particles, as discussed earlier. In addition, Fig. 11 also shows two border upper-cloud scenes denoted by the open triangles at 05:47 and 15:47 UTC, which are among the 26 cases shown in Figs. 10e and 10f. For these border upper-cloud scenes, the mean ARSCL upper CTHs obtained in the

±1.5-min period are lower than 4.3 km and much lower than the MCO2AT and SCO2AT inferred CTHs.

Lastly, Figure 12 shows the histograms of the CTH differences for the comparisons obtained in Figs. 10a-10d. For the overcast upper-cloud scenes (Fig. 12a), the SCO2AT minus ARSCL difference is within ± 0.5 km in 21% of the cases and within ± 1.5 km for 48% of the comparisons. For the MCO2AT cases, the differences within ± 0.5 and ± 1.5 km are improved to 31% and 69%, respectively. Figure 12b shows the comparisons for broken upper-cloud scenes. In the SCO2AT minus ARSCL cases, about 20% are within ± 3 km, which is improved to more than 50% for the MCO2AT cases.

It is clear from these results that the MCO2AT retrieves a value of CTH that is closer to the physical top of the cloud than that obtained with most other passive retrieval methods. Traditionally, the retrieved value of CTH has been associated with the altitude corresponding to the effective radiating temperature of the cloud. For optically thick clouds, that value corresponds more closely to the height at approximately 1 optical depth below the physical cloud top. For optically thin clouds, the retrieved depth is associated with the height of a smaller optical depth. Although some empirical methods have been developed to account for this difference using the window techniques [Minnis et al., 2008, 2009], more evaluation is needed to employ them reliably. By directly retrieving the physical height of the cloud top, the MCO2AT may obviate the need for such corrections.

5. Concluding Remarks

The difficulties of inferring upper cloud top heights (CTHs) for optically thin clouds using passive meteorological satellite data are well known. For upper transmissive clouds, the

inferences become more difficult in multilayer cloud situations when lower clouds coexist underneath the upper clouds. Operational meteorological satellite analyses have shown much success in using the CO₂-slicing or CO₂-absorption techniques to retrieve the CTH information for transmissive cirrus clouds, but all current retrieval algorithms are based on single-layer cloud assumptions. That is, they commonly assume that the cloud layer occupies an infinitesimal thickness in vertical extent and that underneath the geometrically thin cloud layer is a cloud-free atmosphere over the surface.

The single-layer assumptions may induce large uncertainties in the satellite-inferred CTHs because cloud vertical profiles, as revealed by the data obtained from the cloud radar and lidar, are often complex and a single-layer cloud with a geometrically thin thickness is relatively infrequent. To overcome the difficulties with conventional single-layer assumptions, a new method for inferring the upper CTH using a modified CO₂-absorption technique (MCO2AT) has been presented in this paper. This new method employs a three-step sequential retrieval procedure that starts with the single-layer CO₂-absorption technique (SCO2AT) followed by iterative retrieval processes to obtain an enhanced upper CTH. The modified method described in this paper was applied to GOES-12 imager 10.7- and 13.3-µm data and is applicable to data from instruments on other satellites having similar spectral bands. These would include the Spinning Enhanced Visible and Infrared Imager (SEVIRI) onboard Meteosat-8 and -9 and the Moderate-resolution Imaging Spectroradiometer (MODIS) on *Terra* and *Aqua*.

Comparisons between the MCO2AT- and SCO2AT-inferred CTHs and the ARM ground-based ARSCL data demonstrated that the MCO2AT was more effective when the upper portion of a cloud is optically thin but geometrically thick and when an upper cirrus

cloud overlaps with some lower stratus or cumulus clouds. The MCO2AT CTH retrievals represent a significant improvement over the SCO2AT CTH values. Based on comparisons using one month of half-hourly GOES-12 retrievals and the ARSCL data, the mean biases in the GOES-12-inferred upper CTHs were reduced from -1.6 km (SCO2AT) to -0.3 km (MCO2AT) for overcast upper-cloud scenes and from -4.7 km (SCO2AT) to -2.7 km (MCO2AT) for broken upper-cloud scenes. Since the MCO2AT retrieved more accurate upper CTHs, the overall distributions of the upper cloud amounts were enhanced at higher altitudes and the associated root-mean-square errors were also reduced by 30% relative to the SCO2AT retrievals. However, more than 10% of the thin cirrus cloud cases detected by ARSCL were not retrieved by MCO2AT and SCO2AT. This suggests that a more sensitive method is needed to retrieve these thin cirrus clouds. Nevertheless, the comparisons obtained here have provided a preliminary validation for the MCO2AT-inferred CTHs and have shown better agreement with ground-based radar and lidar data, especially for multilayer clouds and optically thin but geometrically thick clouds. Future research plans include applications to other satellite data and additional comparisons with surface and space-based active remote sensing measurements to determine the accuracy of the method over a wide range of cloud, surface, and weather conditions.

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638	Figure Captions
639	Fig. 1 Schematic diagrams of cirrus-overlying-stratus cloud (a, d), geometrically thick and
640	optically thin cloud (b, e), and convective cloud (c, f). P_c from SCO2AT shown on the
641	left (a-c) and P_c and P_{ebg} from MCO2AT are shown on the right (d-f). P_g is ground
642	level.
643	Fig. 2 GOES-12 imager-observed a) 13.3-μm and b) 10.7-μm brightness temperatures, 1 May
644	2005, 10:45 UTC for an area of 32°N-42°N and 105°W-91°W. Square near the center
645	of the images marks the ARM SGP CART site.
646	Fig. 3 a) SCO2AT-inferred P_c , b) ground-level P_g , c) MCO2AT-inferred P_c , and d)
647	MCO2AT-inferred P_{ebg} . Results are obtained for the GOES-12 imager data shown in
648	Fig. 2.
649	Fig. 4 Vertical profiles of MMCR reflectivity (dBZ) obtained at the SGP CART site on 1
650	May 2005. Red arrow indicates scan time of the GOES-12 imager data in Figs. 2 and 3.
651	Fig. 5 Comparisons of a) MCO2AT-inferred P_c and SCO2AT-inferred P_c and b) MCO2AT-
652	inferred P_{ebg} and P_{g} .
653	Fig. 6 GOES-12 imager-observed a) 13.3-μm and b) 10.7-μm brightness temperatures
654	obtained on 24 May 2005 10:45 UTC for an area of 32°N-42°N and 105°W-91°W.
655	Square near the center of the images marks the ARM SGP CART site.
656	Fig. 7 a) SCO2AT-inferred P_c , b) ground-level P_g , c) MCO2AT-inferred P_c , and d)
657	MCO2AT-inferred P_{ebg} . Results are obtained for the GOES-12 imager data shown in
658	Fig. 6.

659	Fig. 8 Vertical profiles of MMCR reflectivity (dBZ) at ARM SGP CART site, 24 May 2005.
660	Red arrow indicates scan time of the GOES-12 imager data in Figs. 6 and 7.
661	Fig. 9 MCO2AT-inferred CTHs (solid points) and SCO2AT-inferred CTHs (cross points)
662	obtained on a) 1 May and b) 24 May 2005. Solid circles indicate overcast upper-cloud
663	scenes and solid triangles indicate broken upper-cloud scenes. Grey areas are cloud
664	masks constructed using the ARSCL data. Arrows indicate the time as shown in Figs.
665	4 and 8.
666	Fig. 10 Comparisons of ARSCL upper CTHs and the MCO2AT (a, c, e) and SCO2AT (b, d, f)
667	inferred CTHs for overcast upper clouds (a, b), broken upper clouds (c, d), and border
668	upper clouds (e, f). Vertical bar indicates the 25th-75th percentiles of the ARSCL data.
669	Fig. 11 CTH retrieved using MCO2AT (solid and open circles) and SCO2AT (crosses) for 26
670	May 2005. Solid circles indicate overcast upper-cloud scenes, solid triangles indicate
671	broken upper-cloud scenes, and open triangles indicate border upper-cloud scenes.
672	Grey areas are cloud masks constructed using the ARSCL data.
673	Fig. 12 Histograms of the upper CTH differences between MCO2AT minus ARSCL (filled
674	grey) and SCO2AT minus ARSCL (no fill). a) for overcast upper-cloud scenes and b)
675	for broken upper-cloud scenes.
676	

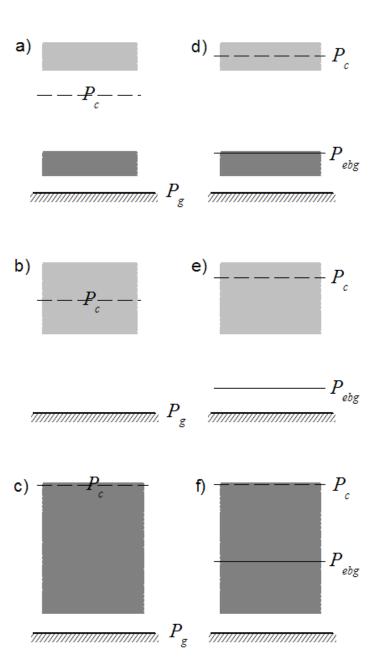


Fig. 1 Schematic diagrams of cirrus-overlying-stratus cloud (a, d), geometrically thick and optically thin cloud (b, e), and convective cloud (c, f). P_c from SCO2AT shown on the left (a-c) and P_c and P_{ebg} from MCO2AT are shown on the right (d-f). P_g is ground level.

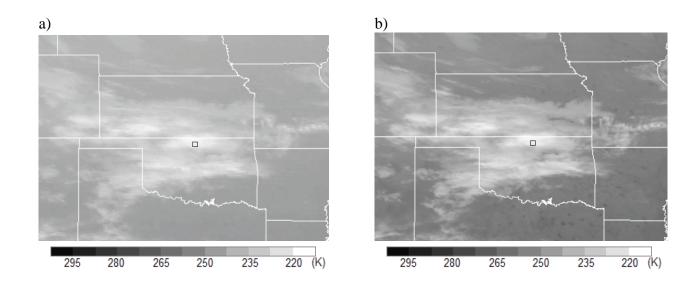


Fig. 2 GOES-12 imager-observed a) 13.3-μm and b) 10.7-μm brightness temperatures, 1 May 2005, 10:45 UTC for an area of 32°N-42°N and 105°W-91°W. Square near the center of the images marks the ARM SGP CART site

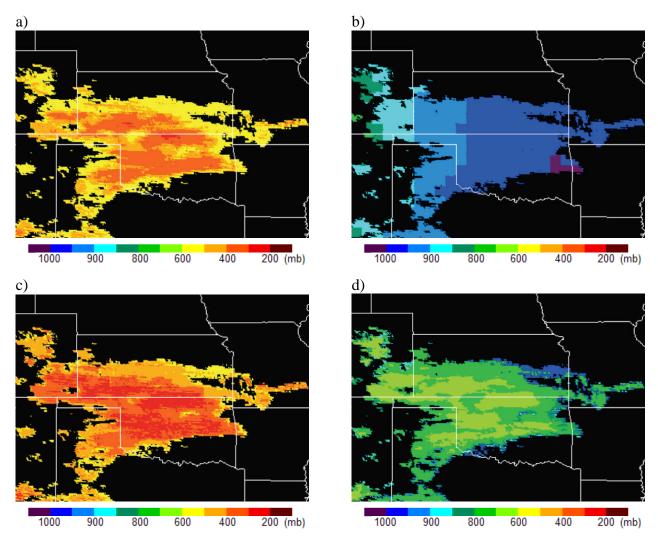


Fig. 3 a) SCO2AT-inferred P_c , b) ground-level P_g , c) MCO2AT-inferred P_c , and d) MCO2AT-inferred P_{ebg} . Results are obtained for the GOES-12 imager data shown in Fig. 2.

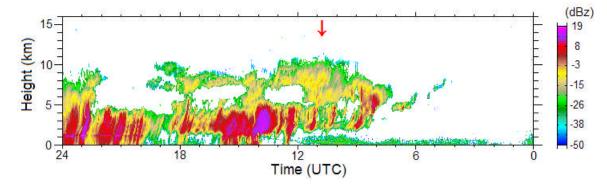


Fig. 4 Vertical profiles of MMCR reflectivity (dBZ) obtained at the SGP CART site on 1 May 2005. Red arrow indicates scan time of the GOES-12 imager data in Figs. 2 and 3.

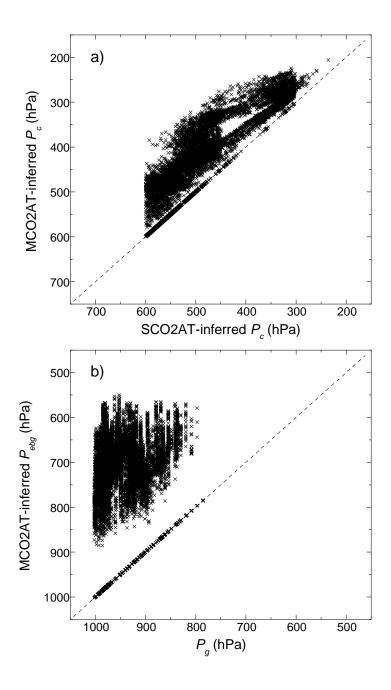


Fig. 5 Comparisons of a) MCO2AT-inferred P_c and SCO2AT-inferred P_c and b) MCO2AT-inferred P_{ebg} and P_g .

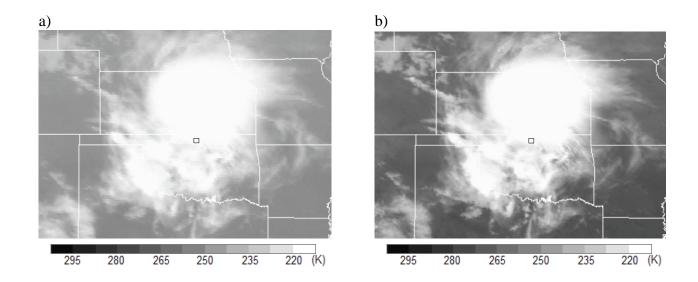


Fig. 6 GOES-12 imager-observed a) 13.3-µm and b) 10.7-µm brightness temperatures obtained on 24 May 2005 10:45 UTC for an area of 32°N-42°N and 105°W-91°W. Square near the center of the images marks the ARM SGP CART site.

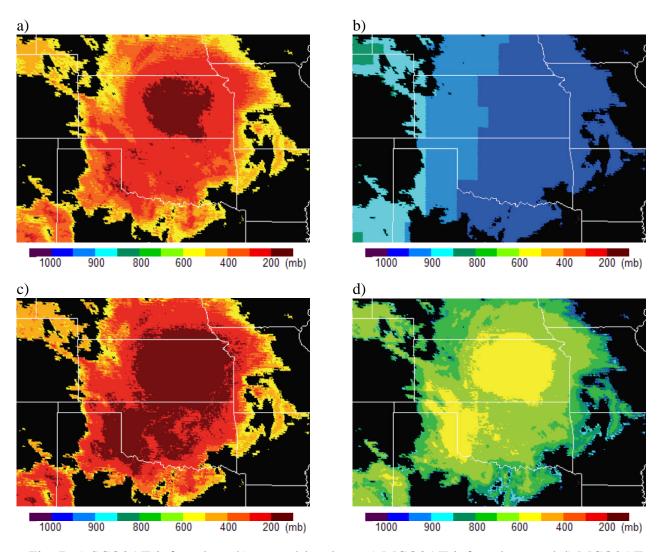


Fig. 7 a) SCO2AT-inferred P_c , b) ground-level P_g , c) MCO2AT-inferred P_c , and d) MCO2AT-inferred P_{ebg} . Results are obtained for the GOES-12 imager data shown in Fig. 6.

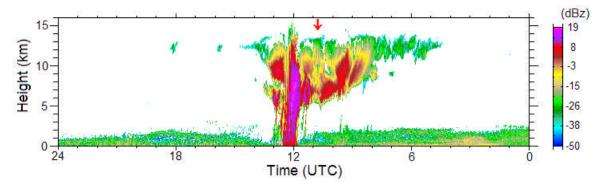


Fig. 8 Vertical profiles of MMCR reflectivity (dBZ) at ARM SGP CART site, 24 May 2005. Red arrow indicates scan time of the GOES-12 imager data in Figs. 6 and 7.

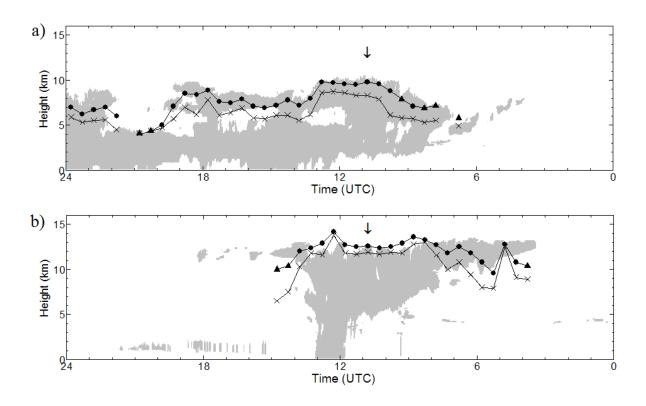


Fig. 9 MCO2AT-inferred CTHs (solid points) and SCO2AT-inferred CTHs (cross points) obtained on a) 1 May and b) 24 May 2005. Solid circles indicate overcast upper-cloud scenes and solid triangles indicate broken upper-cloud scenes. Grey areas are cloud masks constructed using the ARSCL data. Arrows indicate the time as shown in Figs. 4 and 8.

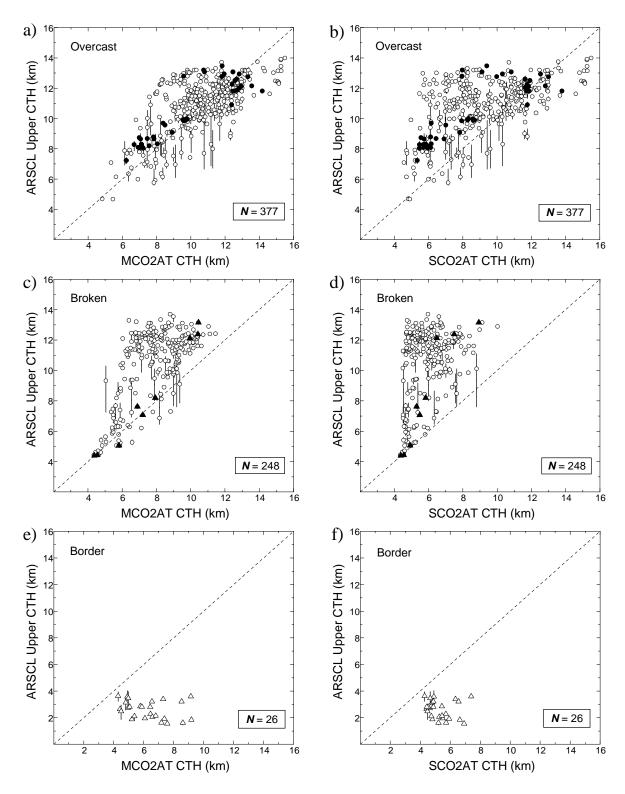


Fig. 10 Comparisons of ARSCL upper CTHs and the MCO2AT (a, c, e) and SCO2AT (b, d, f) inferred CTHs for overcast upper clouds (a, b), broken upper clouds (c, d), and border upper clouds (e, f). Vertical bar indicates the 25th-75th percentiles of the ARSCL data.

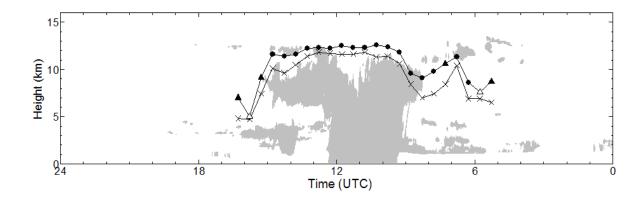


Fig. 11 CTH retrieved using MCO2AT (solid and open circles) and SCO2AT (crosses) for 26 May 2005. Solid circles indicate overcast upper-cloud scenes, solid triangles indicate broken upper-cloud scenes, and open triangles indicate border upper-cloud scenes. Grey areas are cloud masks constructed using the ARSCL data.

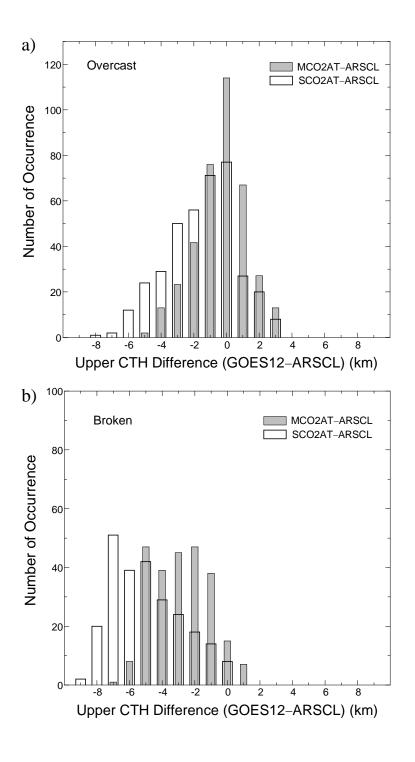


Fig. 12 Histograms of the upper CTH differences between MCO2AT minus ARSCL (filled grey) and SCO2AT minus ARSCL (no fill). a) for overcast upper-cloud scenes and b) for broken upper-cloud scenes.